

Flow-dependent Ekman theory and its application to shallow water models

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Air-sea momentum flux: modified by ocean **surface currents**

1. Wind stress bulk formula

$$\boldsymbol{\tau} = \rho_a c_d |\mathbf{u}_a - \mathbf{u}_o| (\mathbf{u}_a - \mathbf{u}_o)$$

Drag coefficient c_d also depends on ocean surface velocities.

2. Ekman layer dynamics

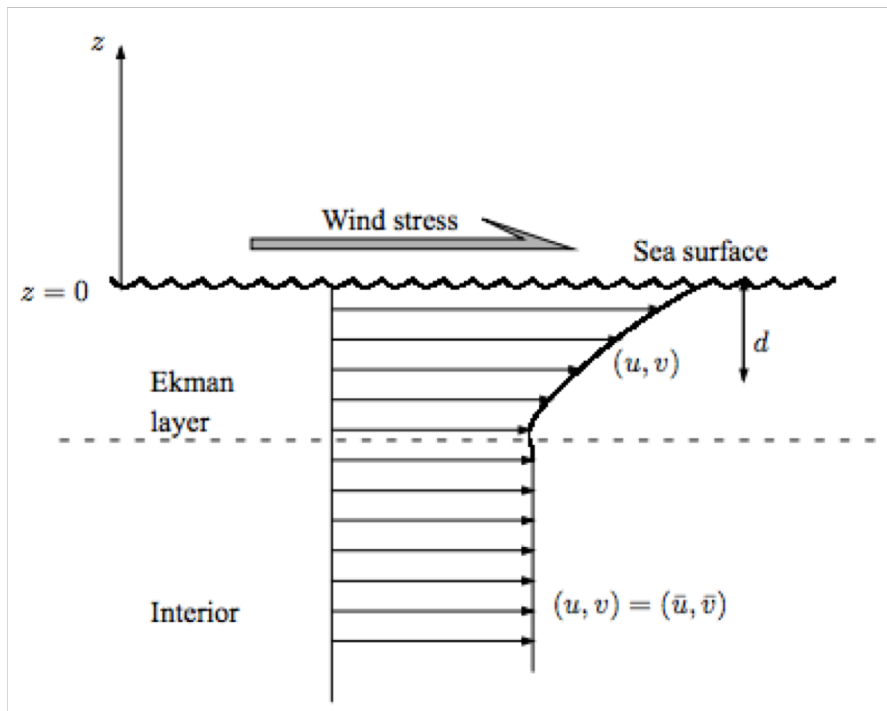
Classic Ekman-layer regime:

$$\begin{cases} f \hat{k} \times \mathbf{u} = \frac{1}{\rho_0} \nabla p + A_z \frac{\partial^2 \mathbf{u}}{\partial z^2} \\ \mathbf{u} = \mathbf{u}_g + \mathbf{u}_{Ek} \end{cases}$$

$$f \hat{k} \times (\mathbf{u}_g + \mathbf{u}_{Ek}) = \frac{1}{\rho_0} \nabla p + A_z \frac{\partial^2 (\mathbf{u}_g)}{\partial z^2} + A_z \frac{\partial^2 (\mathbf{u}_{Ek})}{\partial z^2}$$

Geostrophic balance
(interior ocean)

Ekman balance (surface layer)



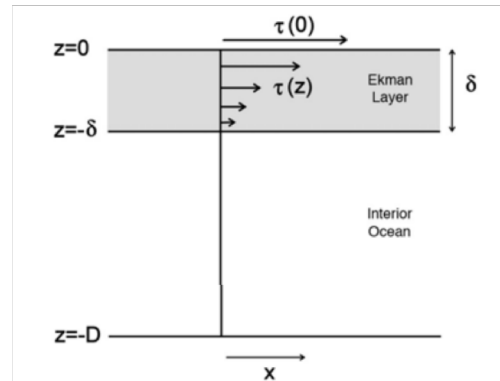
But observations depart significantly from this simple theory.

Development of the flow-dependent Ekman theory

People	Ekman (1905)	Stern and Niiler (1960s)	Wenegrat and Thomas (2017)
Content	Horizontal transport depends on the stress and Coriolis parameter f only.	Allows for shear of the surface velocity field to affect the transport.	Extends early results to better account for curvature in the surface flow path.
Ekman Transport	$U_{Ek} = \frac{\tau_y}{f}$ $V_{Ek} = -\frac{\tau_x}{f}$	$U_{Ek} \approx \frac{\tau_y}{f + \zeta}$ $V_{Ek} \approx -\frac{\tau_x}{f + \zeta}$	$R_0 \bar{u} \frac{\partial V_{Ek}}{\partial s} + (1 + R_0 2\Omega) U_{Ek} = \tau_n$ $R_0 \bar{u} \frac{\partial U_{Ek}}{\partial s} - (1 + R_0 \zeta) V_{Ek} = \tau_s$
Assumptions	Homogeneous deep stationary ocean.	Valid for plane parallel flows (e.g., straight jets); however, not explicitly solved for flows with curvature.	Curvilinear flows, with $R_{oe} \ll 1$ and $R_o < 1$; however, not easily applicable to complicated flow fields.

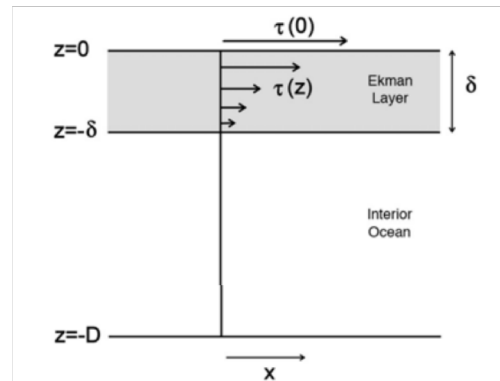
Outline of the following content

1. Flow-dependent Ekman formulation
2. What is the impact of various Ekman formulations on the Ekman transport for **fixed wind stress and oceanic balanced vortex**?



Steady wind stress at the surface;
Constant vortex in the ocean interior.

3. What is the impact of various Ekman formulations on the interior flow when the Ekman layer is **coupled** to the interior?



Seeking solutions for the coupling between
Ekman layer and the interior.

Flow-dependent Ekman layer

$$\frac{\partial \mathbf{u}_{Ek}}{\partial t} + (\mathbf{u}_{Ek} \cdot \nabla) \mathbf{u}_0 + (\mathbf{u}_0 \cdot \nabla) \mathbf{u}_{Ek} + (\mathbf{u}_{Ek} \cdot \nabla) \mathbf{u}_{Ek} + f \hat{k} \times \mathbf{u}_{Ek} = \frac{\partial \boldsymbol{\tau}}{\partial z} - A_h \nabla^4 \mathbf{u}_{Ek} \dots$$

advection 1

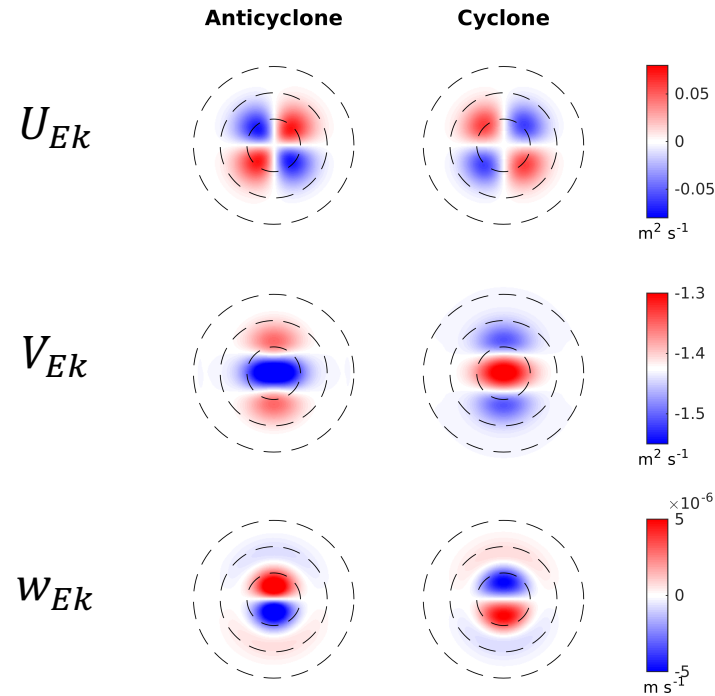
advection 2

nonlinear
(advection 3)

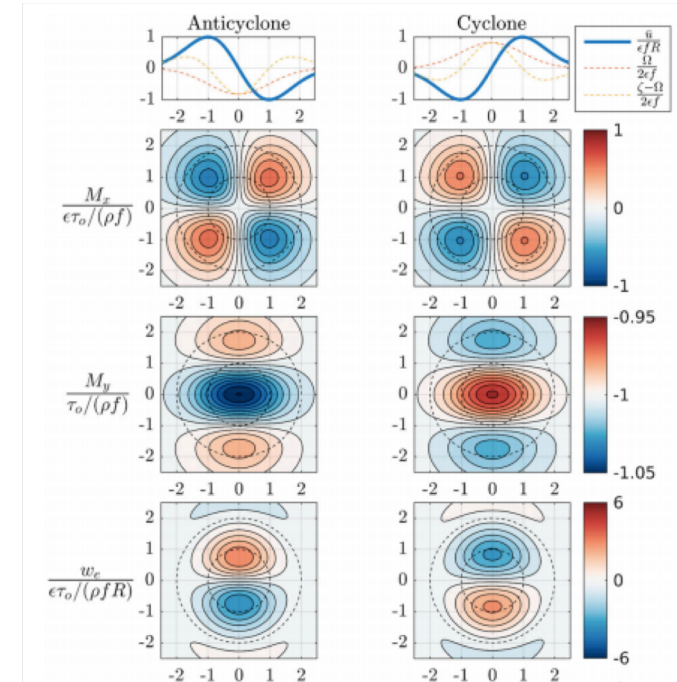
1. The time-dependent term has been added for simplifying the calculation.
2. Scale analysis of the three advection terms: R_o , R_o^2 and $R_o \cdot R_{oe}$.
(Assumptions used here: $R_{oe} \ll 1$ and $R_o < 1$.)
3. Advection 1 has been widely used (or added) to study wind forcing of near-inertial oscillations.
4. Vertical integration leads to the transport equation.

Section 1: the Ekman layer itself

Our Model Simulations



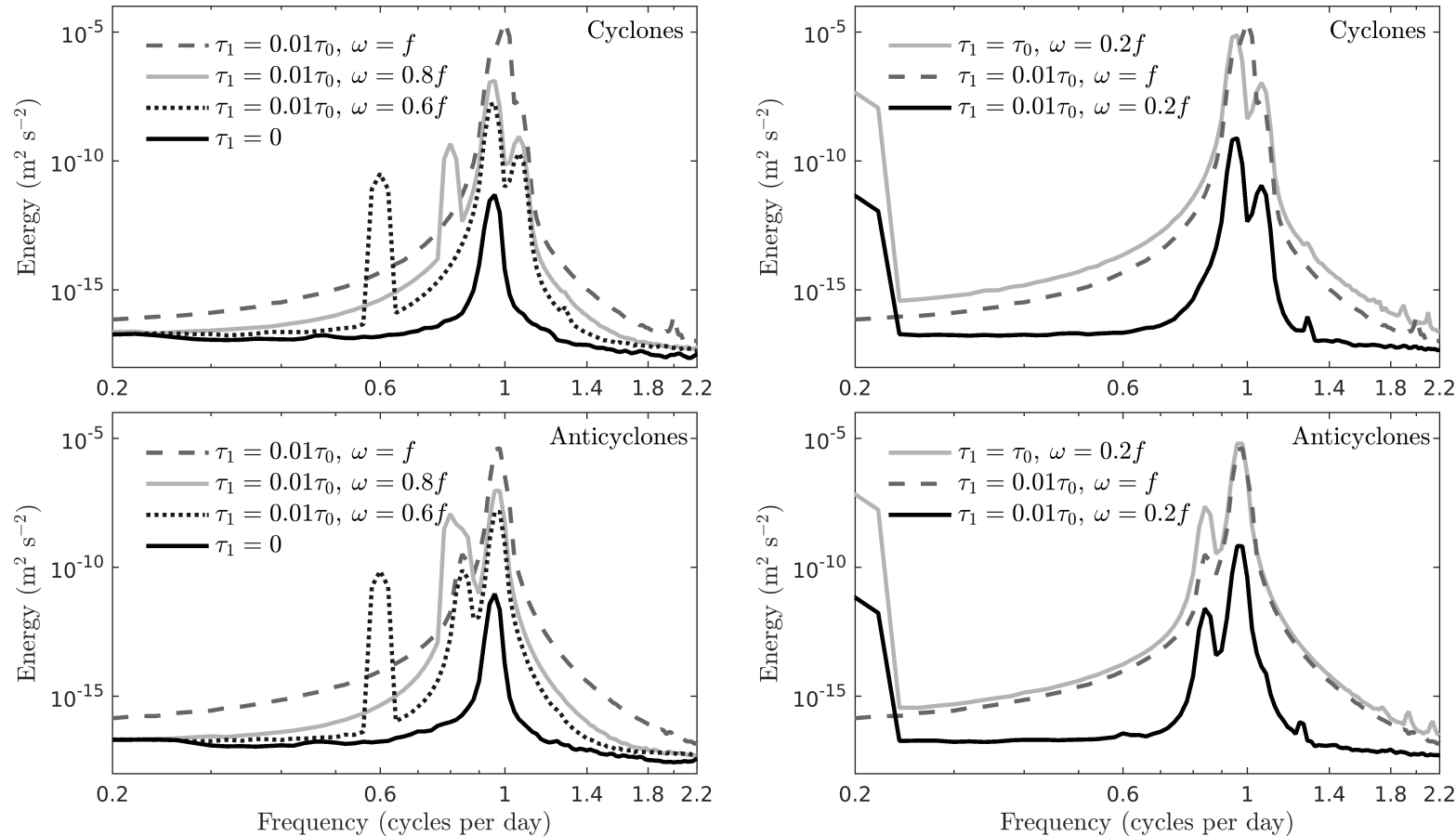
Wenegrat & Thomas simulations



The zonal transport develops a **quadrupole** pattern, emphasizing that the flow-dependent Ekman transport is not strictly perpendicular to the wind stress.

The meridional transport **converges (diverges)** on the north (south) side of the cyclonic vortex, with the pattern reversed for the vortex with anticyclonic flow.

Section 1: the Ekman layer itself



$$\tau_{total} = \tau_0 + \tau_1 \cdot \sin(\omega t)$$

- High-frequency winds lead to responses at the same frequency, plus a component at f .
- Synoptic scale winds with large enough amplitude can be a forcing at f .

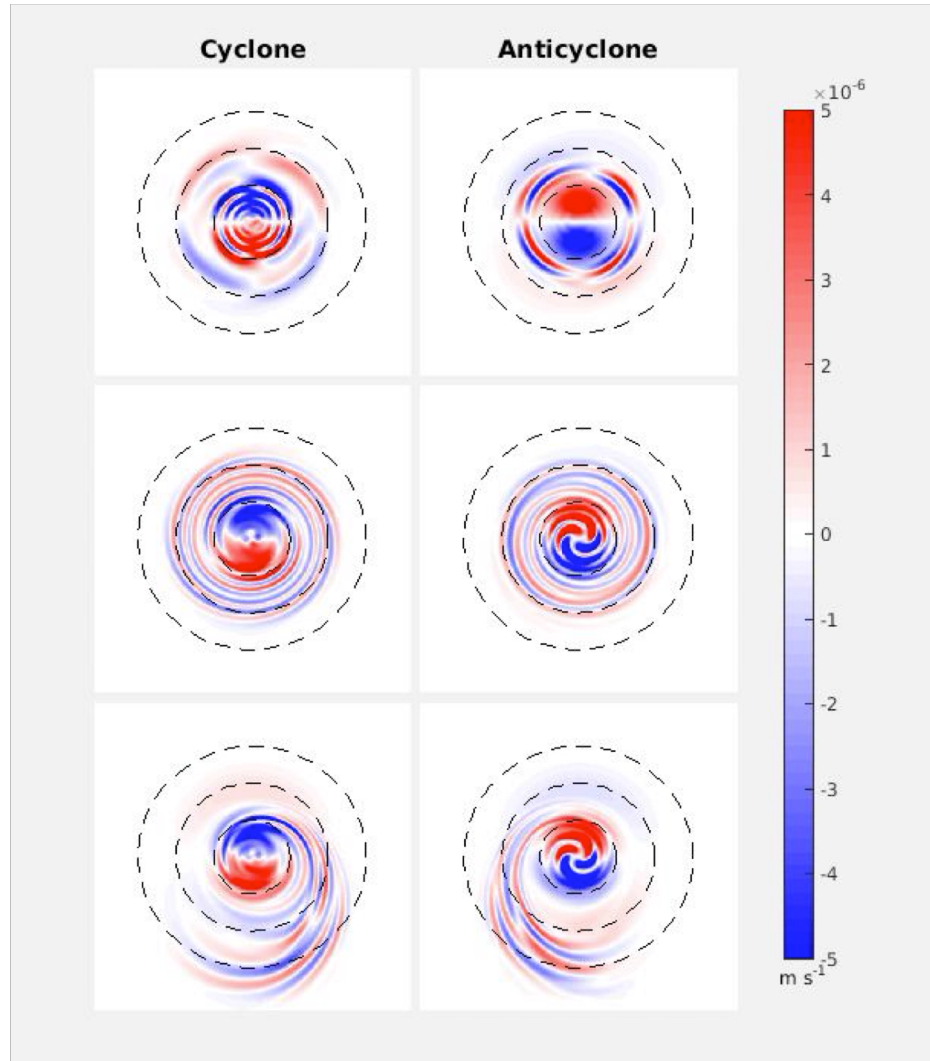
Frequency spectra of pumping velocities with forcing at different frequencies

Section 1: the Ekman layer itself

Advection 1

Advection 1 +
Advection 2

Advection 1 +
Advection 2 +
Nonlinear



Ekman pumping response with different regimes

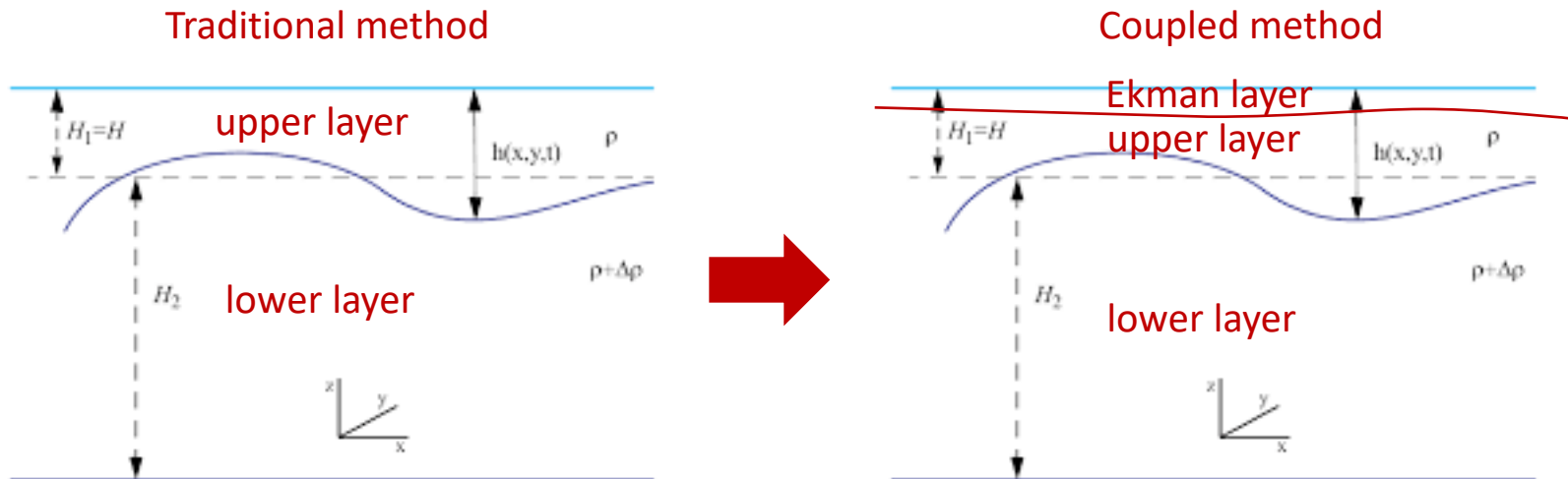
Recap:

- Ekman transport can include a component that is not perpendicular to the stress;
- The time dependence can introduce a near-inertial (high-frequency) component to the pumping velocities.

Section 2: Ekman-interior coupled model

Two different regimes

1. Wind stress is applied as a body force in the upper-layer momentum equation (traditional)
2. Use an explicit Ekman layer to force the upper-layer mass equation (coupled)



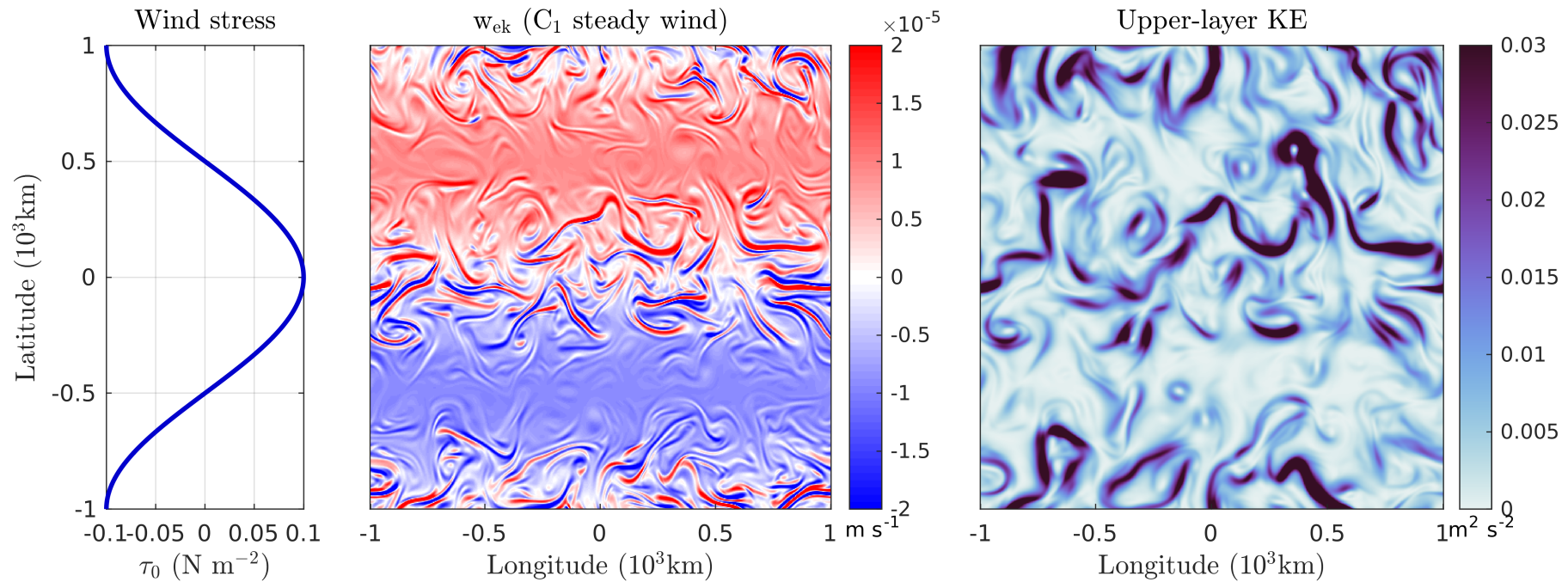
We consider a two-layer shallow water model with a slab Ekman layer in the top layer. Thus, we can use “Ekman pumping” as a forcing in the upper layer mass equation.

Model setup: two-layer rigid lid, domain size (2000km×2000km), resolution (512 grid points×512 grid points), wind forcing τ is a cosine function of latitude.

Section 2: Ekman-interior coupled model

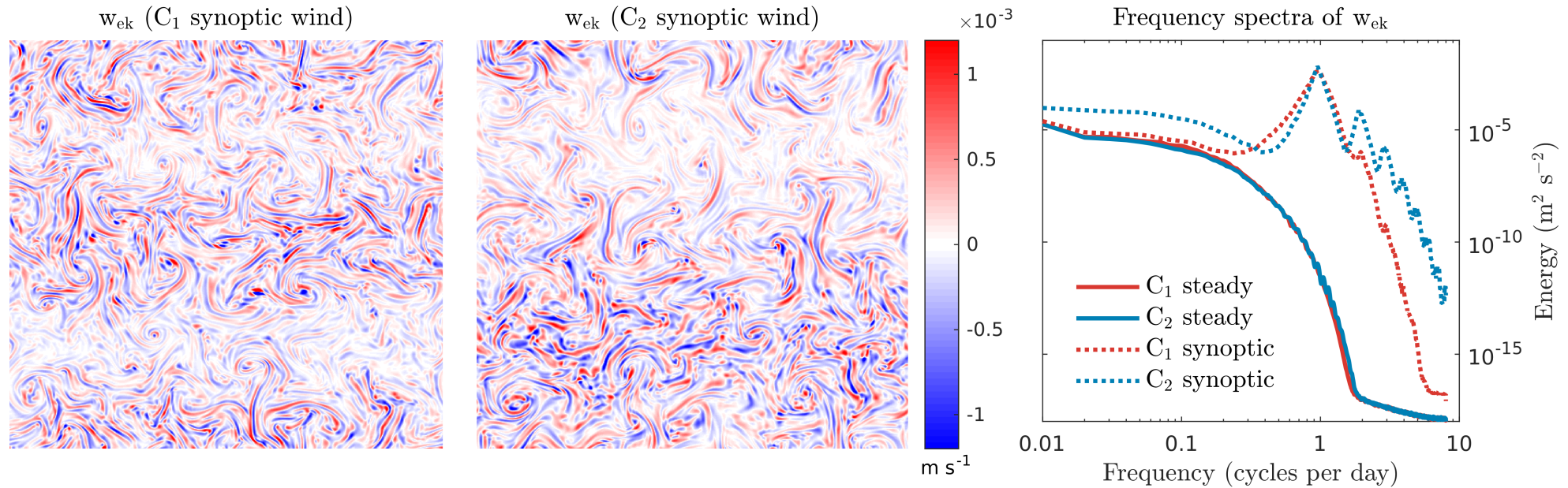
Simulations		Traditional method	Coupled method
Processes		Wind forcing → upper layer	Wind forcing → modified Ekman layer → upper layer
Equations	Ekman layer		Many options as described in model formulation (C1 model: advection1+2; C2 model: advection1+2+3)
	Upper-layer momentum	$\frac{\partial}{\partial t} \vec{u}_1 + (\vec{u}_1 \cdot \nabla) \vec{u}_1 + f \hat{z} \times \vec{u}_1 = \frac{\vec{\tau}}{h_1} - A_h \nabla^4 \vec{u}_1$	$\frac{\partial}{\partial t} \vec{u}_1 + (\vec{u}_1 \cdot \nabla) \vec{u}_1 + f \hat{z} \times \vec{u}_1 = -A_h \nabla^4 \vec{u}_1$
	Upper-layer mass	$\frac{\partial}{\partial t} h_1 + \nabla \cdot (h_1 \vec{u}_1) = 0$	$\frac{\partial}{\partial t} h_1 + \nabla \cdot (h_1 \vec{u}_1) = -w_E$ ($w_E = \nabla \cdot (\vec{U}_E)$)

Section 2: Ekman-interior coupled model



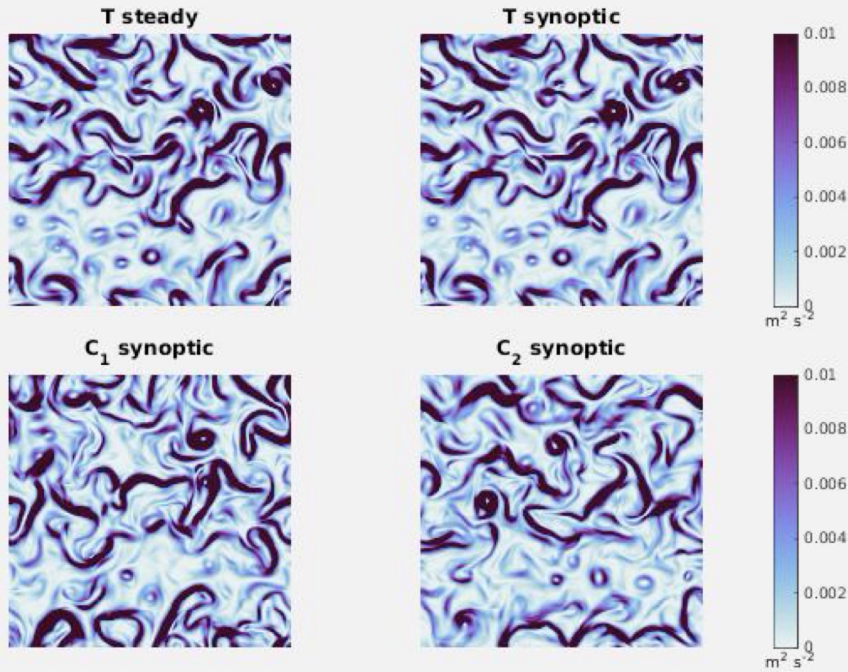
C1 simulation with steady wind
(left: wind structure; mid: Ekman pumping; right: the upper-layer kinetic energy)

Section 2: Ekman-interior coupled model

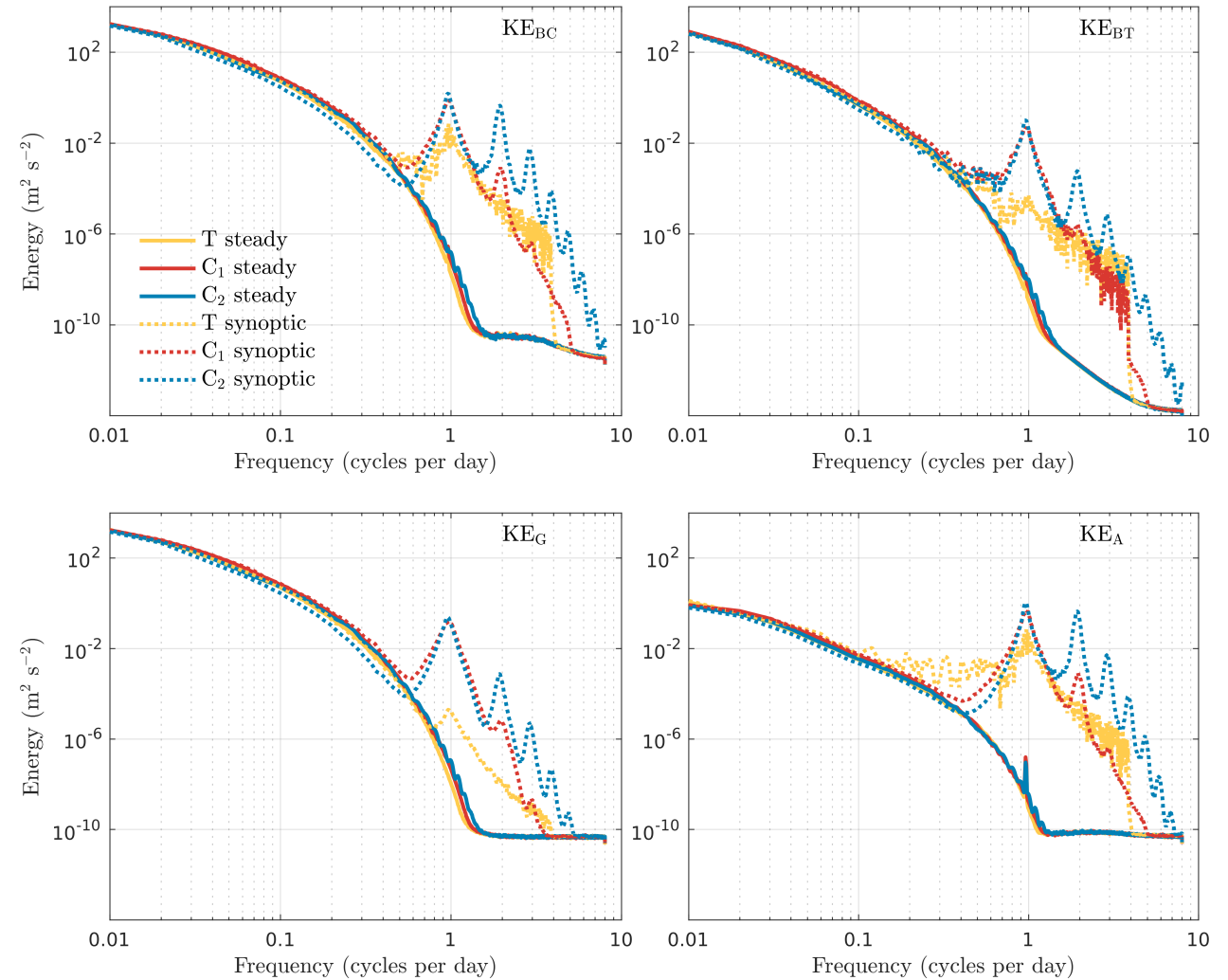


Ekman pumping
(left: C_1 synoptic wind; mid: C_2 synoptic wind; right: frequency spectra)

Section 2: Ekman-interior coupled model



Baroclinic KE of different simulations



Frequency spectra of KE response

Conclusions and discussions

- Flow-dependent Ekman layer can result in a transport that is **not perpendicular** to the wind.
- Synoptic wind can be a **near-inertial forcing** for the flow-dependent Ekman layer.
- With steady wind stress, the Ekman-interior coupled model is almost **identical** to the traditional two-layer shallow water model.
- For the coupled model, adding near-inertial components to the wind stress greatly enhances QG and AG kinetic energy response at **high frequencies**.

Thank you for your attention.